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Role of Magnetic Field on Natural Convective Towards a Semi-Infinite Vertically Inclined Plate in Presence of Hall Current with Numerical Solutions: A Finite Difference Technique

Bandham Saidulu

Assistant Professor, Princeton Degree & PG College, Ramanthapur, Hyderabad, Telangana State, India.

***Corresponding Author:** Bandham Saidulu, Assistant Professor, Princeton Degree & PG College, Ramanthapur, Hyderabad, Telangana State, India.

Abstract: In the present research work, the simultaneous effects of magnetic field and Hall current on an unsteady natural convective flow of incompressible, electrically conducting fluid towards a semi-infinite vertically inclined plate in presence of heat and mass transfer studied. For the numerical results, the basic governing non-linear partial differential equations converted into linear coupled partial differential equations using non-dimensional quintiles and are solved, numerically by finite difference method for velocity (both primary and secondary), temperature, concentration fields are discussed through graphs with the help of engineering parameters such as Grashof number for heat transfer, Grashof number for mass transfer, Prandtl Number, Schmidt Number, Hartmann number and Hall parameter.

 U_{a}

Reference velocity, *m/s*

Keywords: Natural Convection; MHD; Hall current; Finite difference method;

Nomenclature:

List of Variables:

		\mathbf{c}_{o}	1.010101000 (010010),
T'	Temperature of the fluid , K	n_e	Number density of the electron, kg/m^3
C'	Concentration of the fluid, Kg/m^3	n e	rumber density of the election, kg/m
	-		
C'_w	Concentration near the plate, Kg/m^3	M	Hartmann number
		P_{e}	Electron Pressure, N/m^2
C'_{∞}	Concentration in the fluid far away from	-	
	-	е	Electron charge, coulombs
	te, Kg/m^3	Sc	Schmidt Number
T'_w	Temperature of the plate, <i>K</i>		
u'	Velocity component in x' – direction, m/s	Pr	Prandtl number
~	verocity component in <i>x</i> uncertain, <i>m</i> is		
w'	Valasita company in -' direction w/s	Gc	Grashof number for mass transfer
W	Velocity component in z' –direction, m/s		
		Gr	Grashof number for heat transfer
y'	Spatial co-ordinate normal to the plate, m	8	Acceleration due to Gravity, 9.81 m/s^2
x'	Spatial co-ordinate along the plate m	Greek	Symbols:
x'	Spatial co-ordinate along the plate, <i>m</i>	2	Symbols: Volumetric co-efficient of thermal
T'_{∞}	Temperature of the fluid far away from	β	Volumetric co-efficient of thermal
	Temperature of the fluid far away from	etaExpan	Volumetric co-efficient of thermal sion, K^{I}
T'_{∞} the plat	Temperature of the fluid far away from the, K	β	Volumetric co-efficient of thermal
T'_{∞} the plat \overline{V}	Temperature of the fluid far away from te, K Velocity vector, m/s	etaExpan	Volumetric co-efficient of thermal sion, K^{I}
T'_{∞} the plat	Temperature of the fluid far away from the, K	$egin{array}{c} eta \ Expans\ \mu \ lpha \end{array}$	Volumetric co-efficient of thermal sion, K^{-1} Viscosity, <i>Ns/m</i> ² Angle of inclination
T'_{∞} the plat \overline{V} k	Temperature of the fluid far away from te, <i>K</i> Velocity vector, <i>m/s</i> Thermal conductivity, <i>W/mK</i>	etaExpans μ lpha arepsilon	Volumetric co-efficient of thermal sion, K^{-1} Viscosity, <i>Ns/m</i> ² Angle of inclination Porosity of the porous medium
T'_{∞} the plat \overline{V}	Temperature of the fluid far away from te, K Velocity vector, m/s	$egin{array}{c} eta \ Expans\ \mu \ lpha \end{array}$	Volumetric co-efficient of thermal sion, K^{-1} Viscosity, <i>Ns/m</i> ² Angle of inclination
T'_{∞} the plat \overline{V} k	Temperature of the fluid far away from te, <i>K</i> Velocity vector, <i>m/s</i> Thermal conductivity, <i>W/mK</i>	etaExpans μ lpha arepsilon	Volumetric co-efficient of thermal sion, K^{-1} Viscosity, <i>Ns/m</i> ² Angle of inclination Porosity of the porous medium Magnetic permeability, <i>Henry/meter</i>
T'_{∞} the plat \overline{V} k C_p	Temperature of the fluid far away from te, <i>K</i> Velocity vector, <i>m/s</i> Thermal conductivity, <i>W/mK</i> Specific heat at constant Pressure, <i>J/kg-K</i>	etaExpans μ lpha arepsilon	Volumetric co-efficient of thermal sion, K^{-1} Viscosity, <i>Ns/m</i> ² Angle of inclination Porosity of the porous medium
T'_{∞} the plat \overline{V} k	Temperature of the fluid far away from te, <i>K</i> Velocity vector, <i>m/s</i> Thermal conductivity, <i>W/mK</i>	$egin{array}{c} eta & \ { m Expanse} \ \mu & \ lpha & \ arepsilon & \ arepsil$	Volumetric co-efficient of thermal sion, K^{-1} Viscosity, <i>Ns/m</i> ² Angle of inclination Porosity of the porous medium Magnetic permeability, <i>Henry/meter</i>
T'_{∞} the plat \overline{V} k C_p	Temperature of the fluid far away from te, <i>K</i> Velocity vector, <i>m/s</i> Thermal conductivity, <i>W/mK</i> Specific heat at constant Pressure, <i>J/kg-K</i>	$eta \ {\cal B} \ { m Expans} \ { m \mu} \ { m lpha} \ { m eta} \ {$	Volumetric co-efficient of thermal sion, K^{l} Viscosity, <i>Ns/m</i> ² Angle of inclination Porosity of the porous medium Magnetic permeability, <i>Henry/meter</i> Kinematics viscosity, m^{2}/s
T'_{∞} the plat \overline{V} k C_p k_e	Temperature of the fluid far away from te, <i>K</i> Velocity vector, <i>m/s</i> Thermal conductivity, <i>W/mK</i> Specific heat at constant Pressure, <i>J/kg-K</i> Mean absorption coefficient	$egin{array}{c} eta & \ { m Expanse} \ \mu & \ lpha & \ arepsilon & \ arepsil$	Volumetric co-efficient of thermal sion, K^{-1} Viscosity, <i>Ns/m</i> ² Angle of inclination Porosity of the porous medium Magnetic permeability, <i>Henry/meter</i>
T'_{∞} the plat \overline{V} k C_p	Temperature of the fluid far away from te, <i>K</i> Velocity vector, <i>m/s</i> Thermal conductivity, <i>W/mK</i> Specific heat at constant Pressure, <i>J/kg-K</i>	$eta \ {\cal B} \ { m Expans} \ { m \mu} \ { m lpha} \ { m eta} \ {$	Volumetric co-efficient of thermal sion, K^{l} Viscosity, <i>Ns/m</i> ² Angle of inclination Porosity of the porous medium Magnetic permeability, <i>Henry/meter</i> Kinematics viscosity, m^{2}/s

$ au_{e}$	Electron collision time, Sec	β^* Co-efficient of volume expansion with
σ	Electrical conductivity, <i>mho/m</i>	Species concentration, m^3/Kg Superscripts:
θ	Dimensionless Temperature (K)	Dimensionless Properties Subscripts:
ϕ	Dimensionless concentration (Kg/m^3)	W Wall condition
ρ	Density, kg/m^3	p Plate
,	••••	∞ Free stream condition

1. INTRODUCTION

The studies of free convective flow of viscous incompressible fluid past vertical bodies have wide range of extensive engineering and technological applications. When free convection flows occurs at high temperature, the effects of radiation are found to be posing vital important applications. Thermal radiation is key to many fundamental phenomena surrounding atmosphere, from solar radiation to fire and radiant lamp, which had played a major role in combustion and furnace design, design of fins, nuclear power plants, cooling of towers, gas turbines and various propulsion devices for aircraft, energy utilization, temperature measurements, remote sensing for astronomy, and space exploration. Magneto hydrodynamic flow of heat and mass transfer processes occur in many of the industrial applications: such as cooling of geothermal systems, aerodynamic processes, chemical catalytic reactors and processes, electromagnetic pumps, and Magneto Hydrodynamic power generators etc. Many studies have been carried out to investigate the magneto hydrodynamic past free convective fluid flow. The study of magnetohydrodynamic flow through porous media has been the subject of considerable improvement in the last few decades. In nuclear engineering, the magneto hydrodynamic flow in a porous medium is required for the design of a layer of liquid metal around a thermonuclear fusion-fission hybrid reactor. In metallurgy, a permanent magnetic field can be applied during the solidification process to modify the intensity of the interdendritic flow of the metallic liquid in a porous medium. Dharmendar Reddy et al. ([1] and [2]) studied hall current effect on an unsteady MHD free convection flow past a vertical porous plate with chemical reaction, heat and mass transfer. Srinivasa Rao [3] studied finite element analysis of radiation and mass transfer flow past semi-infinite moving vertical plate with viscous dissipation. Sheri et al. ([4]-[7]) discussed transient approach to heat absorption and radiative heat transfer past an impulsively moving plate with ramped temperature. Venkataramana et al. [8] studied thermal radiation and rotation effect on an unsteady MHD mixed convection flow through a porous medium with Hall current and Heat absorption. Sailaja et al. ([9]-[12]) discussed finite element solutions of non-newtonian dissipative Casson fluid flow past a vertically inclined surface surrounded by porous medium including constant heat flux, thermal diffusion and diffusion thermo. Sivaiah and Srinivasa Raju and their co-authors ([13]-[14]) studied finite element solution of heat and mass transfer flow with hall current, heat source and viscous dissipation. Ramya et al. ([15]-[18]) discussed boundary layer viscous flow of nanofluids and heat transfer over a nonlinearly isothermal stretching sheet in the presence of heat generation/absorption and slip boundary conditions. Srinivasa Raju ([19] and [20]) studied unsteady MHD boundary layer flow of casson fluid over an inclined surface embedded in a porous medium with thermal radiation and chemical reaction. Srinivasa Raju et al. ([21] and [22]) studied radiation effect on unsteady MHD free convection with Hall current near on an infinite vertical porous plate. Srinivasa Raju [23] studied numerical treatment of casson fluid free convection flow past an infinite vertical plate filled in magnetic field in presence of angle of inclination and thermal radiation using finite element technique. Srinivasa Raju et al. ([24]-[39]) discussed influence of transpiration on unsteady heat transfer MHD fluid flow over an infinite vertical plate in presence of hall current. Srinivasa Raju [40] studied effects of soret and Dufour on natural convective fluid flow past a vertical plate embedded in porous medium in the presence of thermal radiation via finite element method. Srinivasa Raju et al. [41] discussed magnetohydrodynamic chemically reacting flow past vertically inclined permeable plate filled in porous medium with convergence analysis of FEM. Combined influence of thermal diffusion and diffusion thermo on unsteady hydromagnetic free convective fluid flow past an infinite vertical porous plate in presence of chemical reaction studied by Srinivasa Raju [42]. Srinivasa Raju et al. ([43]-[46]) discussed the application of finite element method to unsteady MHD free convection flow past a vertically inclined porous plate including thermal diffusion and diffusion thermo effects. Application of finite element method to MHD mixed convection chemically reacting flow past a vertical porous

plate with cross diffusion and Biot number effects studied by Srinivasa Raju [47]. Simultaneous effects of soret and Dufour on unsteady hydromagnetic free convective chemically reacting fluid flow past an infinite vertical plate filled in porous medium studied by Srinivasa Raju and Sarada [48]. Study of grid independence of finite element method on MHD free convective casson fluid flow with slip effect studied by Srinivasa Raju and Ramesh [49]. Manideep et al. ([50] and [51]) discussed MHD free convection heat transfer Couette flow in rotating system. Maddilety and Srinivasa Raju [52] studied hall effect on an unsteady MHD free convective Couette flow between two permeable plates. Heat and mass transfer effects on MHD natural convective flow past an infinite vertical porous plate with thermal radiation and hall current studied by Ramana Murthy et al. [53]. Sudhakar et al. ([54]-[56]) studied hall effect on an unsteady MHD flow past along a porous flat plate with thermal diffusion, diffusion thermo and chemical reaction. Unsteady MHD free convection flow near on an infinite vertical plate embedded in a porous medium with chemical reaction, hall current and thermal radiation studied by Sarada et al. [57]. Anand Rao et al. ([58]-[67]) studied finite element analysis of unsteady MHD free convection flow past an infinite vertical plate with Soret, Dufour, thermal radiation and heat source. Anand Rao and Srinivasa Raju ([68]-[70]) studied hall effect on an unsteady MHD flow and heat transfer along a porous flat plate with mass transfer and viscous dissipation. Jitthender Reddy et al. ([71]-[79]) discussed chemical reaction and radiation effects on MHD free convection from an impulsively started infinite vertical plate with viscous dissipation. Aruna et al. [80] studied combined influence of Soret and Dufour effects on unsteady hydromagnetic mixed convective flow in an accelerated vertical wavy plate through a porous medium, Krishna Prasad et al. ([81]-[83]) discussed thermal radiation influence on MHD Flow of a rotating fluid with heat transfer through finite element and element free Galerkin solutions.

Motivated by the above reference work, the aim of this investigation is to study the influence of hall current on an unsteady magneto hydrodynamic flow and mass transfer of an electrically conducting incompressible fluid along an infinite vertically inclined plate. Also, the effects of different flow parameters encountered in the equations are studied. The problem is governed by system of coupled non-linear partial differential equations whose exact solution is difficult to obtain. Hence, the problem is solved by using finite difference method, which is more economical from computational view point.

2. MATHEMATICAL FORMULATION

An unsteady natural convection flow of an electrically conducting viscous incompressible fluid with mass transfer along a porous flat plate has been considered with hall current. In Cartesian coordinate system, let x' - axis is taken to be along the plate and the y' - axis normal to the plate. Since the plate is considered infinite in x'-direction, hence all physical quantities will be independent of x'direction. Let the components of velocity along x' and y' axes be u' and y' which are chosen in the upward direction along the plate and normal to the plate respectively. A uniform magnetic field of magnitude B_{o} is applied normal to the plate. The transverse applied magnetic field and magnetic Reynold's number are assumed to be very small, so that the induced magnetic field is negligible. The polarization effects are assumed to be negligible and hence the electric field is also negligible. Initially, for time $t' \leq 0$, the plate and the fluid are maintained at the same constant temperature (T'_{α}) in a stationary condition with the same species concentration (C'_{∞}) at all points so that, the Soret and Dufour effects are neglected. When t' > 0, The wall is maintained at constant temperature (T'_w) and concentration (C'_{w}) higher than the ambient temperature (T'_{w}) and concentration (C'_{w}) respectively. The homogeneous chemical reaction of first order with rate constant between the diffusing species and the fluid is assumed. Using the relation $\nabla \overline{H} = 0$ for the magnetic field $\overline{H} = (H_x, H_y, H_z)$, we obtain $H_y = \text{constant} = H_a$ (say) where H_a is the externally applied transverse magnetic field so that $\overline{H} = (0, H_o, 0)$. The equation of conservation of electric charge $\nabla \cdot \overline{J} = 0$ gives $j_y = \text{constant}$, where $\overline{J} = (j_x, j_y, j_z)$. We further assume that the plate is non-conducting. This implies $j_y = 0$ at the plate and hence zero everywhere. When the strength of magnetic field is very large, the generalized Ohm's law in the absence of electric field takes the following form:

$$\overline{J} + \frac{\omega_e \tau_e}{B_o} \overline{J} \times \overline{H} = \sigma \left(\mu_e \overline{V} \times \overline{H} + \frac{1}{en_e} \nabla P_e \right)$$
(1)

Under the assumption that the electron pressure (for weakly ionized gas), the thermo-electric pressure and ion-slip conditions are negligible, equation (1) becomes:

$$j_{x} = \frac{\sigma \mu_{e} H_{o}}{1 + m^{2}} (mu' - w') \text{ and } j_{z} = \frac{\sigma \mu_{e} H_{o}}{1 + m^{2}} (mw' + u')$$
(2)

Where u' is the x'-component of \overline{V} , w' is the z'-component of \overline{V} and $m (= \omega_e \tau_e)$ is the hall parameter. Within the above framework, the equations which govern the flow under the usual Boussinesq's approximation are as follows:

$$\frac{\partial v'}{\partial y'} = 0 \tag{3}$$

$$\frac{\partial u'}{\partial t'} + v' \frac{\partial u'}{\partial y'} = v \frac{\partial^2 u'}{\partial y'^2} - \frac{\sigma \mu^2 {}_e H_o^2}{\rho (1+m^2)} (u' + mw') + g\beta (T' - T'_{\infty}) (\cos \alpha) + g\beta^* (C' - C'_{\infty}) (\cos \alpha)$$
(4)

$$\frac{\partial w'}{\partial t'} + v' \frac{\partial w'}{\partial y'} = v \frac{\partial^2 w'}{\partial y'^2} - \frac{\sigma \mu^2 {}_e H^2_{o}}{\rho (1+m^2)} (w' - mu')$$
(5)

$$\frac{\partial T'}{\partial t'} + v' \frac{\partial T'}{\partial y'} = \frac{k}{\rho C_p} \frac{\partial^2 T'}{\partial {y'}^2}$$
(6)

$$\frac{\partial C'}{\partial t'} + v' \frac{\partial C'}{\partial y'} = D \frac{\partial^2 C'}{\partial {y'}^2}$$
(7)

The corresponding initial and boundary conditions are:

$$t' \le 0: \quad u' = 0, \quad w' = 0, \quad T' = T'_{\infty}, \quad C' = C'_{\infty} \quad for \; all \quad y'$$

$$t' > 0: \begin{cases} u' = 0, \quad w' = 0, \quad T' = T'_{w}, \quad C' = C'_{w} \quad at \quad y' = 0\\ u' = 0, \quad w' = 0, \quad T' = T'_{\infty}, \quad C' = C'_{\infty} \quad as \quad y' \to \infty \end{cases}$$
(8)

The non-dimensional quantities introduced in the equations (3)-(7) are:

$$t = \frac{t'U_{o}^{2}}{\upsilon}, \quad y = \frac{y'U_{o}^{2}}{\upsilon}, (u, v, w) = \frac{(u', v', w')}{U_{o}}, \quad \theta = \frac{(T' - T'_{w})}{(T'_{w} - T'_{w})}, \quad \varphi = \frac{(C' - C'_{w})}{(C'_{w} - C'_{w})}, \quad M = \frac{\sigma\mu_{e}^{2}H_{0}^{2}\upsilon}{\rho U_{o}^{2}}, \quad Sc = \frac{\upsilon}{D}, \quad Gr = \frac{\upsilon g\beta(T'_{w} - T'_{w})}{U_{o}^{3}}, \quad Gr = \frac{\upsilon g\beta(T'_{w} - T'_{w})}{U_{o}^{3}}, \quad Gr = \frac{\omega g\beta(T'_{w} - T'_{w})}{U_{o}^{3}}, \quad Gr = \frac{\omega$$

The governing equations can be obtained in the dimensionless form as:

$$\frac{\partial v}{\partial y} = 0 \tag{10}$$

$$\frac{\partial u}{\partial t} + v \frac{\partial u}{\partial y} = \frac{\partial^2 u}{\partial y^2} - \frac{M}{\left(1 + m^2\right)} \left(u + mw\right) + (Gr)\theta\left(\cos\alpha\right) + (Gc)\phi\left(\cos\alpha\right)$$
(11)

$$\frac{\partial w}{\partial t} + v \frac{\partial w}{\partial y} = \frac{\partial^2 w}{\partial y^2} - \frac{M}{\left(1 + m^2\right)} \left(w - mu\right) \tag{12}$$

$$\frac{\partial \theta}{\partial t} + v \frac{\partial \theta}{\partial y} = \frac{1}{\Pr} \frac{\partial^2 \theta}{\partial y^2}$$
(13)

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$$\frac{\partial \phi}{\partial t} + v \frac{\partial \phi}{\partial y} = \frac{1}{Sc} \frac{\partial^2 \phi}{\partial y^2}$$
(14)

The initial and boundary conditions (8) in the non-dimensional form are:

$$t \le 0: u = 0, \ w = 0, \ \theta = 0, \ \phi = 0 \quad for all \quad y \\ t > 0: \begin{cases} u = 0, \ w = 0, \ \theta = 1, \ \phi = 1 & at \ y = 0 \\ u = 0, \ w = 0, \ \theta = 0, \ \phi = 0 & as \ y \to \infty \end{cases}$$
(15)

3. RESULTS AND DISCUSSIONS

Hall effect study on an unsteady magnetohydrodynamic flow and mass transfer of an electrically conducting incompressible newtonian fluid along an infinite vertically inclined plate with heat and mass transfer has been studied and solved by using finite difference method. The effects of material parameters such as Grashof number for heat transfer, Grashof number for mass transfer, Prandtl number, Schmidt number, Hartmann number and Hall parameter separately in order to clearly observe their respective effects on the primary velocity, secondary velocity, temperature and concentration profiles of the flow. During the course of numerical calculations of the primary velocity, secondary velocity, temperature and concentration the values of the Prandtl number are chosen for Mercury (Pr = 0.025), Air at $25^{\circ}C$ and one atmospheric pressure (Pr = 0.71), Water (Pr = 7.00) and Methanol (Pr = 11.62). To focus out attention on numerical values of the results obtained in the study, the values of Sc are chosen for the gases representing diffusing chemical species of most common interest in air namely Hydrogen (Sc = 0.22), Helium (Sc = 0.30), Water-vapour (Sc = 0.60), Oxygen (Sc = 0.66) and Ammonia (Sc = 0.78). For the physical significance, the numerical discussions in the problem and at t = 1.0, stable values for primary velocity, secondary velocity, temperature and concentration fields are obtained. To examine the effect of parameters related to the problem on the velocity field and skin-friction numerical computations are carried out at Pr = 0.71. To find out the solution of this problem, we have placed an infinite vertical plate in a finite length in the flow. Hence, we solve the entire problem in a finite boundary. However, in the graphs, the y values vary from 0 to 9, and the primary velocity, secondary velocity, temperature and concentration profiles tend to zero as y tends to 9. This is true for any value of y. Thus, we have considered finite length.

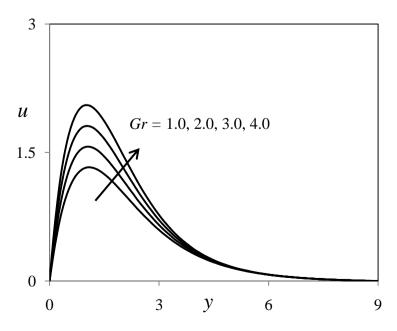


Fig. 1. Influence of Gr on primary velocity profiles

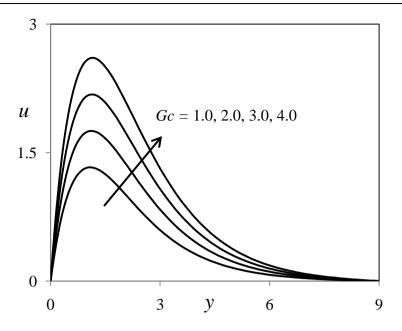


Fig. 2. Influence of Gc on primary velocity profiles

The temperature and the species concentration are coupled to the velocity via Grashof number for heat and mass transfer as seen in equation (11). For various values of Grashof number for heat and mass transfer, the primary velocity profiles u are plotted in figures (1) and (2). The Grashof number for heat transfer signifies the relative effect of the thermal buoyancy force to the viscous hydrodynamic force in the boundary layer. As expected, it is observed that there is a rise in the primary velocity due to the enhancement of thermal buoyancy force. Also, as Gr increases, the peak values of the primary velocity increases rapidly near the porous plate and then decays smoothly to the free stream velocity. The Grashof number for mass transfer defines the ratio of the species buoyancy force to the viscous hydrodynamic force. As expected, the fluid velocity increases and the peak value is more distinctive due to increase in the species buoyancy force. The primary velocity distribution attains a distinctive maximum value in the vicinity of the plate and then decreases properly to approach the free stream value. It is noticed that the velocity increases with increasing values of Modified Grashof number. The nature of primary velocity profiles in presence of foreign species such as Hydrogen (Sc = 0.22), Helium (Sc = 0.30), Oxygen (Sc = 0.60) and Ammonia (Sc = 0.78) are shown in figure (3). The flow field suffers a decrease in primary velocity at all points in presence of heavier diffusing species.

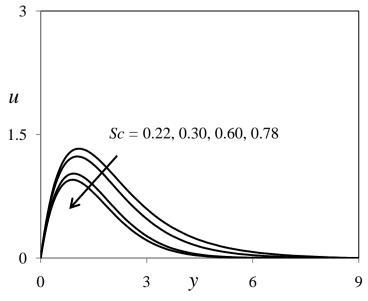


Fig. 3. Influence of Sc on primary velocity profiles

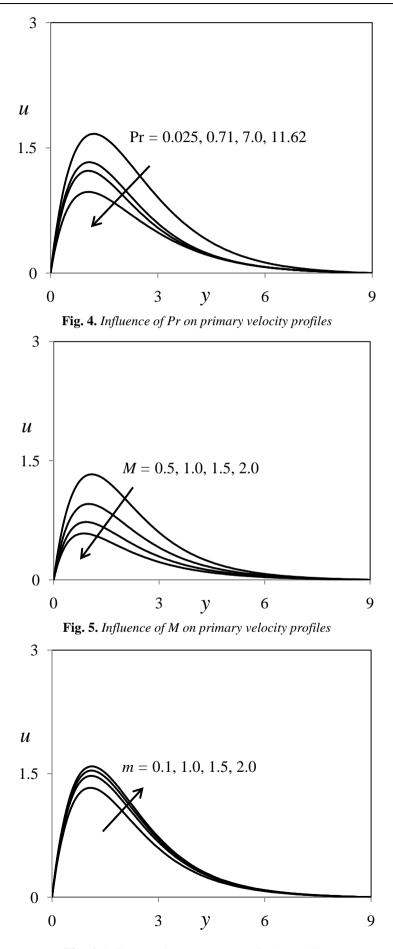


Fig. 6. Influence of m on primary velocity profiles

The influence of the viscous dissipation parameter i.e., the Eckert number on the velocity and temperature are shown in figures (7) and (8) respectively. The Eckert number expresses the relationship between the kinetic energy in the flow and the enthalpy. It embodies the conversion of kinetic energy into internal energy by work done against the viscous fluid stresses. Greater viscous dissipative heat causes a rise in the temperature as well as the velocity. This behavior is evident from figures (7) and (8). In figures (9) and (10), we see the influence of the both heat and mass transfer on secondary velocity of the flow. It can be seen that as both the heat and mass transfer increases, this velocity component increases as well. In figures (9) and (10), the effects of the heat and mass transfer on the velocity are displayed. It is apparent from the figures that, the increasing values of heat and mass transfer enhance the secondary velocity. The effect of Eckert number on the secondary velocity flow field is presented in the figure (11). Here, the secondary velocity profiles are drawn against yfor three different values of Ec. The Eckert number is found to accelerate the secondary velocity of the flow field at all points. In figure (12) we have the influence of the Hartmann number on the secondary velocity. It can be seen that as the values of this parameter increases, the secondary velocity increases. The nature of secondary velocity profiles in presence of foreign species such as Hydrogen (Sc = 0.22), Helium (Sc = 0.30), Water-vapour (Sc = 0.60) and Ammonia (Sc = 0.78) are shown in figure (13). The flow field suffers a decrease in secondary velocity at all points in presence of heavier diffusing species. Figure (4) depicts the effect of Prandtl number on primary velocity profiles in presence of foreign species such as Mercury (Pr = 0.025), Air at 25° C and one atmospheric pressure (Pr = 0.71), Water (Pr = 7.00) and Methanol (Pr = 11.62) are shown in figure (4). We observe that from figure (4), the primary velocity decreases with increasing of Prandtl number. The effect of the Hartmann number is shown in figure (5). It is observed that, the primary velocity of the fluid decreases with the increasing of the magnetic field number values. The decrease in the primary velocity as the Hartmann number increases is because the presence of a magnetic field in an electrically conducting fluid introduces a force called the Lorentz force, which acts against the flow, if the magnetic field is applied in the normal direction, as in the present study. This resistive force slows down the fluid velocity component as shown in figure (5). Figure (6) depicts the primary velocity profiles as the Hall parameter m increases. We see that u increases as m increases. It can also be observed that u profiles approach their classical values when Hall parameter m becomes large

(m > 1).

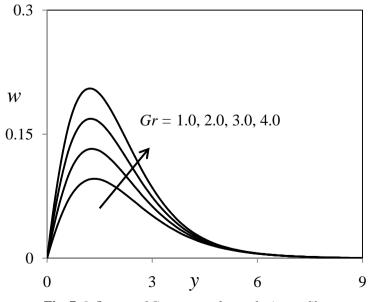


Fig. 7. Influence of Gr on secondary velocity profiles

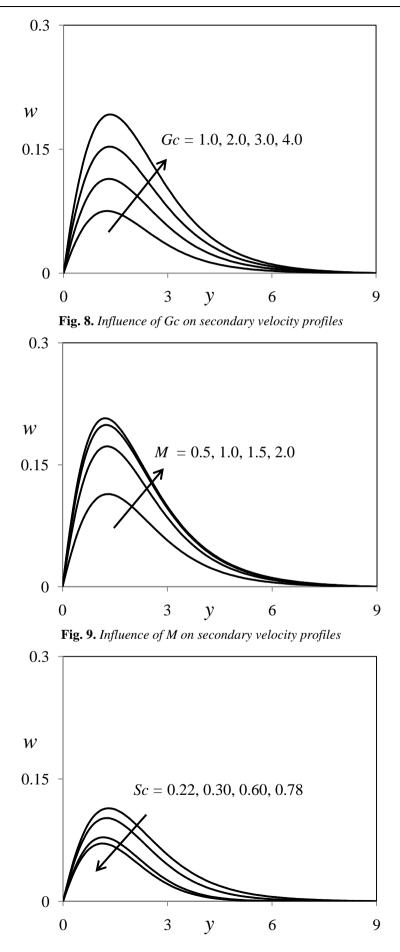


Fig. 10. Influence of Sc on secondary velocity profiles

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In figures (7) and (8), we see the influence of the both heat and mass transfer on secondary velocity of the flow. It can be seen that as both the heat and mass transfer increases, this velocity component increases as well. In figures (7) and (8), the effects of the heat and mass transfer on the velocity are displayed. It is apparent from the figures that, the increasing values of heat and mass transfer enhance the secondary velocity. In figure (9), the influence of the Hartmann number on the secondary velocity. It can be seen that as the values of this parameter increases, the secondary velocity increases. The nature of secondary velocity profiles in presence of foreign species such as Hydrogen (Sc = 0.22), Helium (Sc = 0.30), Water-vapour (Sc = 0.60) and Ammonia (Sc = 0.78) are shown in figure (10). The flow field suffers a decrease in secondary velocity at all points in presence of heavier diffusing species. Figure (11) depicts the effect of Prandtl number on secondary velocity profiles in presence of foreign species such as Mercury (Pr = 0.025), Air at $25^{\circ}C$ and one atmospheric pressure (Pr = 0.71), Water (Pr = 7.00) and Methanol (Pr = 11.62) are shown in figure (11). It is observed that, the velocity is decreasing with increasing of Prandtl number. In figure (12) we depict the effect of Prandtl number on the temperature field. It is observed that an increase in the Prandtl number leads to decrease in the temperature field. Also, temperature field falls more rapidly for water in comparison to air and the temperature curve is exactly linear for mercury, which is more sensible towards change in temperature. From this observation it is concluded that mercury is most effective for maintaining temperature differences and can be used efficiently in the laboratory. Air can replace mercury, the effectiveness of maintaining temperature changes are much less than mercury. However, air can be better and cheap replacement for industrial purpose. This is because, either increase of kinematic viscosity or decrease of thermal conductivity leads to increase in the value of Prandtl number. Hence temperature decreases with increasing of Prandtl number. The effects of Schmidt number on the concentration field is presented in figure (13). Figure (13) shows the concentration field due to variation in Schmidt number for the gasses Hydrogen, Helium, Water-vapour, Oxygen and Ammonia. It is observed that, the concentration field is steadily for Hydrogen and falls rapidly for Oxygen and Ammonia in comparison to Water-vapour. Thus Hydrogen can be used for maintaining effective concentration field and Water-vapour can be used for maintaining normal concentration field.

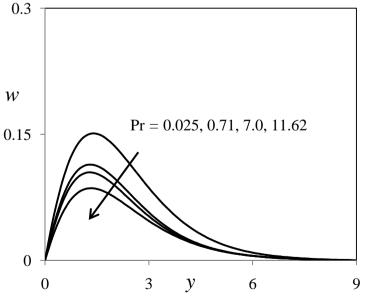


Fig. 11. Influence of Pr on secondary velocity profiles

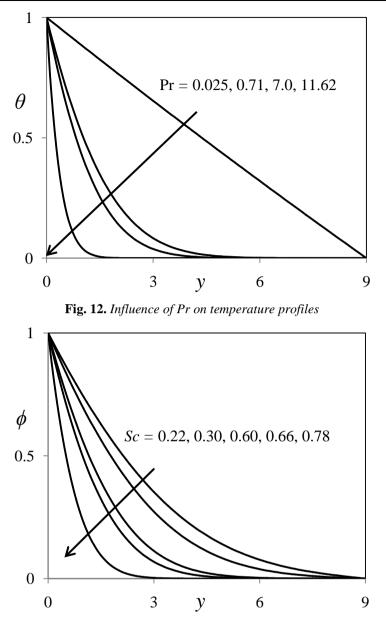


Fig. 13. Influence of Sc on concentration profiles

4. CONCLUSIONS

The effect of hall current on unsteady MHD flow of an electrically conducting incompressible fluid along a porous flat plate with heat and mass transfer is studied. The dimensionless equations are solved by using finite difference method. The simultaneous effects of primary velocity, secondary velocity, temperature and concentration for different parameters like Grashof number for heat transfer, Grashof number for mass transfer, Prandtl number, Schmidt number, Hartmann number and Hall parameter studied. The study concludes the following results.

- 1. It is observed that both the primary and secondary velocities of the fluid increases with the increasing of parameters Grashof number for heat transfer, Grashof number for mass transfer, Hall parameter and decreases with the increasing of parameters Prandtl number, Schmidt number, Hartmann number.
- 2. The fluid temperature increases with the decreasing of Prandtl number.
- 3. The Concentration of the fluid decreases with the increasing of Schmidt number.

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