

Non-Vacuum LRS Bianchi Type-V Cosmological Model in $f(R)$ Theory of Gravity

L.S. Ladke¹

Department of Mathematics
Nilkanthrao Shinde Science and Arts college
Bhandravati, India
lemrajvandana@gmail.com

R. A. Hiwarkar²

Department of Mathematics
Guru Nanak Institute of Engg. & Technology
Nagpur, India
rhiwarkar@gmail.com

V. K. Jaiswal³

Department of Mathematics
Priyadarshini J. L. College of Engineering,
Nagpur, India
vkjassi@gmail.com

Abstract: *We study the non-vacuum perfect fluid solution of LRS Bianchi type-V space-time in the $f(R)$ theory of gravity to obtain energy density and pressure of the universe by using stiff matter. The time dependent deceleration parameter is used to discuss the physical behavior of the model. The function of the Ricci scalar is also evaluated.*

Keywords: *LRS Bianchi type-V space-time, $f(R)$ theory of gravity, deceleration parameter.*

1. INTRODUCTION

The general theory of relativity has the tremendous success which explains the most of gravitational phenomena, but it fails to explain the nature of the negative component known as dark energy and accelerating expansion of the universe. Therefore there is a need of modification to general theory of relativity. Recent observations also strongly support that the universe is an accelerated state which is conformed from CMB [1-6]. Plenty of modifications had been done. Among them each one has its novel features. The $f(R)$ theory of gravity is one of the modified theory which answered the problem of dark matter, dark energy and accelerating expansion of the universe. This theory has such important characteristic, hence it is considered as most suitable.

Many authors explained the concept of $f(R)$ theory of gravity. The late time acceleration and early time deceleration confirmed by Nijiri and Odinstov [7] in $f(R)$ theory of gravity. The Nojiri and Odintsov [8] have studied various modified gravity theories that are considered as gravitational alternatives for dark energy. M. Sharif and H. Rizwana Kausar [9] explained anisotropic Bianchi type-III model in $f(R)$ gravity by assuming exponential and power-law volumetric expansion which represent an accelerated expansion of the universe. M. Sharif et. al. [10] explained the energy distribution in $f(R)$ gravity by using generalized Landau-Lifshitz energy-momentum complex. K. S. Adhav [11] obtained Bianchi type-III string cosmological model in $f(R)$ gravity. Hollenstein and Francisco S. N. Lobo [12] discussed the exact solutions of $f(R)$ gravity coupled to nonlinear electrodynamics. Many authors have done wonderful work on $f(R)$ theory of gravity in different context.

The investigation of Bianchi type models in modified $f(R)$ theory of gravity is an interesting discussion. The anomalies found in the cosmic microwave background (CMB) and large scale structure observations stimulated a growing interest in anisotropic cosmological models of the universe. Kumar and Singh [13] solved Bianchi type-I space-time in general relativity in the presence of perfect fluid. M. Sharif and H. Rizwana Kausar [14] discussed non-vacuum solutions of Bianchi

Type-VI₀ universe in $f(R)$ gravity considering isotropic perfect fluid. M. Sharif and M. Farasat Shamir [15] exhibited perfect fluid solutions of Bianchi types-I and V space-times in $f(R)$ theory of gravity. The Bianchi models being anisotropic are useful to study isotropic behavior of the universe with the passage of time. Sharif and Kausar [16] studied the isotropic behavior of Bianchi III model in $f(R)$ gravity. The scalar field can play a vital role to explain the cosmic acceleration which has widely been studied in $f(R)$ gravity [17,18]. As $f(R, T)$ theory involves coupling between matter and geometry, so considering scalar field as a source may provide some new insights. Farasat Shamir et. al.[19] obtained exact solutions of Bianchi types-I &V models in $f(R, T)$ gravity which gives constant deceleration parameter by using variation law of Hubble parameter. R.L.Naidu et. al.[20] discussed FRW viscous fluid cosmological model in $f(R, T)$ gravity. H.R. Ghate et. al.[21] studied Bianchi type-IX viscous string cosmological model in $f(R, T)$ gravity with special form of deceleration parameter. P.K. Sahoo et. al. [22] investigated Kaluza-Klein cosmological model in $f(R, T)$ gravity with $\Lambda(t)$ by considering constant decelerating parameter.

Bianchi models play a significant role because these models are homogeneous and anisotropic in nature, in which the process of isotropization of the universe is studied through the passage of time. The study of Bianchi type-V cosmological model create more interest because these models are anisotropic generalization of open FRW model and allow arbitrarily small anisotropy levels at any constant of cosmic time. Bianchi type-V model have been studied in details by number of authors viz. Banerjee and Sanyal [23], Collins [24], Wainwrigth et al. [25].

We are attentive in the $f(R)$ theory of gravity. With such motivation, in this paper, we have examined non-vacuum perfect fluid solution of LRS Bianchi type-V space-time in the $f(R)$ theory of gravity. We used stiff matter to obtained energy density and pressure of the universe. Also we used time dependent deceleration parameter to discuss the physical behavior of the model. The function of the Ricci scalar is also evaluated.

2. FIELD EQUATIONS IN $f(R)$ THEORY OF GRAVITY

The $f(R)$ theory of gravity is the simplest generalization of the general theory of relativity proposed by Einstein in which Ricci scalar in Einstein–Hilbert action is replaced by an arbitrary function of the Ricci scalar. The field equations in $f(R)$ theory of gravity are given by

$$F(R)R_{ij} - \frac{1}{2}f(R)g_{ij} - \nabla_i \nabla_j F(R) + g_{ij} \square F(R) = kT_{ij}, \tag{1}$$

where $F(R) \equiv \frac{df(R)}{dR}$, $\square \equiv \nabla^i \nabla_i$, ∇_i is the covariant derivative and T_{ij} is the standard matter energy momentum tensor.

If we take $f(R) = R$, the field equations (1) reduce to the field equation of general theory of relativity which is proposed by Einstein.

3. LRS BIANCHI TYPE-V MODEL

The line element of LRS Bianchi type-V space-time is given by

$$ds^2 = dt^2 - A^2 dx^2 - e^{2mx} B^2 [dy^2 + dz^2], \tag{2}$$

where A, B are called cosmic scale factors which are function of time t and m is an arbitrary constant.

The energy momentum tensor for perfect fluid has of the form

$$T_{ij} = (\rho + p)u_i u_j - pg_{ij} \tag{3}$$

satisfying the equation of state

$$p = w\rho, \quad 0 \leq w \leq 1, \tag{4}$$

where p is pressure, ρ is the energy density while $u^i = \delta_4^i$ is the four velocity vector which satisfies

$$u^i u_i = 1 . \tag{5}$$

From equation (1), we have the field equations

$$-2 \frac{\ddot{B}}{B} + 2 \frac{\dot{A}\dot{B}}{AB} + \frac{\dot{A}\dot{F}}{AF} - \frac{\ddot{F}}{F} - \frac{2m^2}{A^2} = \frac{k}{F}(\rho + p) . \tag{6}$$

$$-\frac{\ddot{A}}{A} - \frac{\ddot{B}}{B} + \frac{\dot{A}\dot{B}}{AB} + \frac{\dot{B}^2}{B^2} + \frac{\dot{B}\dot{F}}{BF} - \frac{\ddot{F}}{F} - \frac{2m^2}{A^2} = \frac{k}{F}(\rho + p) . \tag{7}$$

$$\left(\frac{\dot{A}}{A} - \frac{\dot{B}}{B} \right) = 0 . \tag{8}$$

On integrating equation (8), we have

$$A = k B \tag{9}$$

where k is constant of integration

The average scale factor & volume scale factor defined as

$$a = (AB^2)^{\frac{1}{3}} \tag{10}$$

$$V = a^3 = AB^2 . \tag{11}$$

We also defined the generalized Hubble parameter H and deceleration parameter q as

$$H = \frac{\dot{a}}{a} = \frac{1}{3} [H_1 + H_2 + H_3] , \tag{12}$$

$$q = -\frac{a\ddot{a}}{\dot{a}^2} = -\frac{\ddot{a}}{aH^2} \tag{13}$$

where $H_1 = \frac{\dot{A}}{A}$, $H_2 = H_3 = \frac{\dot{B}}{B}$ are the directional Hubble parameters in the direction of x, y, z axes respectively.

The expansion scalar θ and shear scalar σ^2 are defined as follows

$$\theta = u^i_{;i} = \frac{\dot{A}}{A} + \frac{2\dot{B}}{B} , \tag{14}$$

$$\sigma^2 = \frac{1}{2} \sigma_{ij} \sigma^{ij} = \frac{1}{3} \left[\frac{\dot{A}}{A} - \frac{\dot{B}}{B} \right]^2 , \tag{15}$$

where
$$\sigma_{ij} = \frac{1}{2} (\nabla_j u_i + \nabla_i u_j) - \frac{1}{3} g_{ij} \theta . \tag{16}$$

The mean anisotropy parameter A_m is given by

$$A_m = \frac{1}{3} \sum_{i=1}^3 \left(\frac{H_i - H}{H} \right)^2 . \tag{17}$$

After solving (6) and (7), we have

$$\frac{A}{B} = d_1 \exp\left[c_1 \int \frac{dt}{a^3 F}\right], \tag{18}$$

where d_1 and c_1 are constants of integration.

Using equation (11) and (18), the metric functions turn out to be

$$A = ap_1 \exp\left[q_1 \int \frac{dt}{a^3 F}\right], \tag{19}$$

$$B = ap_2 \exp\left[q_2 \int \frac{dt}{a^3 F}\right], \tag{20}$$

where $p_1 = d_1^{\frac{2}{3}}, p_2 = d_1^{-\frac{1}{3}},$ (21)

and $q_1 = \frac{2}{3}c_1, q_2 = -\frac{1}{3}c_1,$ (22)

Satisfying the relations

$$p_1 p_2^2 = 1, \quad q_1 + 2q_2 = 0. \tag{23}$$

In order to solve the Einstein’s field equations in $f(R)$ theory of gravity, we used the a special form of deceleration parameters which defined by Abdussatter and Prajapati [26], as

$$q = -\frac{\alpha}{t^2} + \beta - 1, \tag{24}$$

where $\alpha > 0, \beta > 1$ is a constant.

Equation (13) can be integrated to obtain the scale factor as

$$a(t) = e^{k_2} \exp\left[\int \frac{1}{\int (1+q)dt + k_1} dt\right]. \tag{25}$$

where k_1 and k_2 are constants of integration.

Using equations (24) and (25), we calculated mean scale factor by substituting $k_1 = 0$ as

$$a(t) = e^{k_2} \left(t^2 + \frac{\alpha}{\beta}\right)^{\frac{1}{2\beta}}, \tag{26}$$

Sharif and Shamir [15] have established a result in the context of $f(R)$ gravity which show that

$$F \propto a^b. \tag{27}$$

Thus, we have

$$F = l a^b, \tag{28}$$

where l is the constant of proportionality, b is any integer.

We use the value of $a(t)$ for $k_2 = 0, b = -1, \beta = 1$ in the equations (19) and (20), we obtained the scale factor as

$$A = p_1 (t^2 + \alpha)^{\frac{1}{2}} \exp\left(\frac{q_1}{l\sqrt{\alpha}} \tan^{-1} \frac{t}{\sqrt{\alpha}}\right). \tag{29}$$

$$B = p_2 (t^2 + \alpha)^{1/2} \exp\left(\frac{q_2}{l\sqrt{\alpha}} \tan^{-1} \frac{t}{\sqrt{\alpha}}\right), \quad (30)$$

Hence geometry of the universe (2) is given by

$$ds^2 = dt^2 - (t^2 + \alpha) \left\{ p_1^2 \exp\left(\frac{2q_1}{l\sqrt{\alpha}} \tan^{-1} \frac{t}{\sqrt{\alpha}}\right) dx^2 - e^{2mx} p_2^2 \exp\left(\frac{2q_2}{l\sqrt{\alpha}} \tan^{-1} \frac{t}{\sqrt{\alpha}}\right) [dy^2 + dz^2] \right\} \quad (31)$$

The volume scale factor is given by

$$V = (t^2 + \alpha)^{3/2}. \quad (32)$$

The directional Hubble parameters in the directions of x, y and z axis are found to be

$$H_1 = \frac{q_1}{l(t^2 + \alpha)} + \frac{t}{(t^2 + \alpha)}, \quad (33)$$

$$H_2 = H_3 = \frac{q_2}{l(t^2 + \alpha)} + \frac{t}{(t^2 + \alpha)} \quad (34)$$

The mean Hubble parameter becomes

$$H = \frac{t}{(t^2 + \alpha)}, \quad (35)$$

The expansion scalar θ is given by

$$\theta = \frac{3t}{(t^2 + \alpha)}, \quad (36)$$

The shear scalar σ^2 calculated as

$$\sigma^2 = \frac{(q_1 - q_2)^2}{3l^2(t^2 + \alpha)^2}. \quad (37)$$

The mean anisotropy parameter A_m turns out to be

$$A_m = \frac{2(q_1 - q_2)^2}{9l^2 t^2}, \quad (38)$$

The energy density and pressure of the universe for stiff matter obtained as

$$2k\rho = 2kp = \frac{(t^2 + \alpha)^{1/2}}{l} \left[\frac{(q_1 - 6q_2)t}{l} + \frac{2q_2(q_2 - q_1)}{l^2} + t^2 - \alpha \right] - \frac{2m^2(t^2 + \alpha)}{p_1^2} \exp\left(\frac{-2q_1}{l\sqrt{\alpha}} \tan^{-1} \frac{t}{\sqrt{\alpha}}\right), \quad (39)$$

However, function of Ricci scalar for Bianchi type-V model is given by

$$f(R) = \frac{1}{(t^2 + \alpha)^{5/2}} \left[\frac{3m^2 l (t^2 + \alpha)}{p_1^2} \exp\left(\frac{-2q_1}{l\sqrt{\alpha}} \tan^{-1} \frac{t}{\sqrt{\alpha}}\right) - \frac{3l(5t^2 + 3\alpha)}{2} + \frac{q_1(2q_1 + q_2)}{l} \right] + \frac{(t^2 + \alpha)^3}{2l} \left(\frac{(q_1 - 6q_2)t}{l} + \frac{2q_2(q_2 - q_1)}{l^2} + t^2 - \alpha \right), \quad (40)$$

where

$$R = \frac{2}{(t^2 + \alpha)} \left[\frac{q_2(2q_1 + q_2)}{l^2(t^2 + \alpha)} + \frac{3m^2}{p_1^2} \exp\left(\frac{-2q_1}{l\sqrt{\alpha}} \tan^{-1} \frac{t}{\sqrt{\alpha}}\right) - 3 \right]. \quad (41)$$

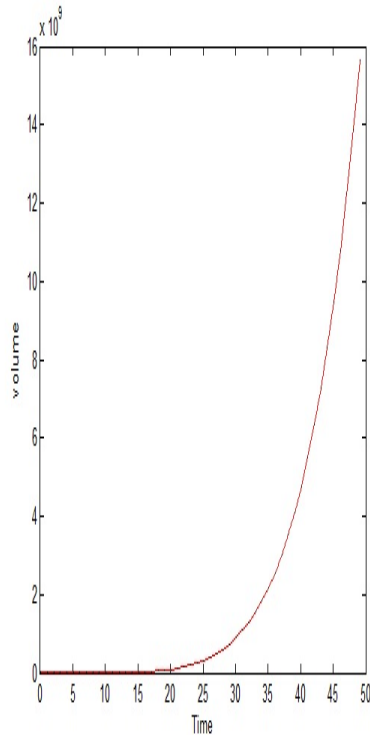


Fig.(1) Volume vs Time

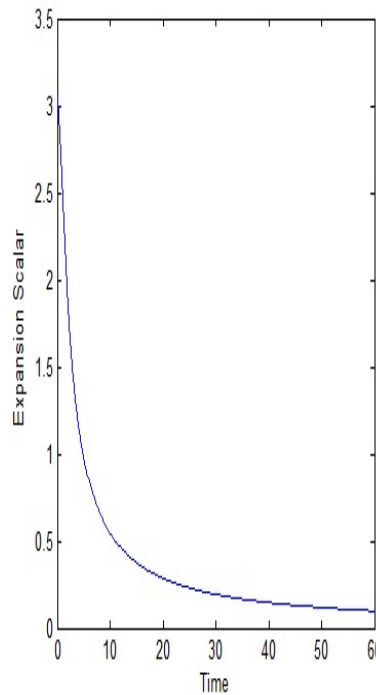


Fig.(2) Expansion Scalar vs Time

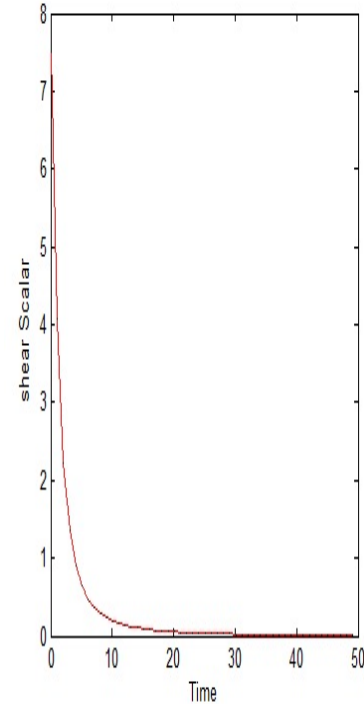


Fig.(3) Shear scalar vs Time

4. CONCLUDING REMARK

In this paper, we have examined the non-vacuum perfect fluid solution of LRS Bianchi type-V space-time in the $f(R)$ theory of gravity. We used stiff matter to obtain energy density and pressure of the universe. Also the time dependent deceleration parameter proposed by Abdussatter and Prajapati helped to discuss the physical behavior of the model. The function of the Ricci scalar is also evaluated. The observation is as under:

- 1) The model is free from initial singularity.
- 2) The spatial volume V is zero at $t \rightarrow 0$ and increases infinitely with time. This shows that at the initial epoch the universe starts with zero volume and expands uniformly as shown in Fig.(1).
- 3) It is interesting to note that the present model is shear free and isotropic throughout expansion of the universe if $q_1 = q_2$, otherwise model is not shear free and it is an anisotropic.
- 4) The graph of expansion scalar and shear scalar versus time indicates that shear scalar & expansion scalar are large as $t=0$ and become constant when time increases as shown in Fig.(2) & Fig.(3) respectively.

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