

## Higher Five Dimensional Plane Gravitational Waves with Wet Dark Energy in Bimetric Relativity

S. D. Deo

Department of Mathematics  
N.S.Science and Arts College,  
Bhadrawati, Dist-Chandrapur, (M.S.)442902,  
India  
s\_deo01@yahoo.in

Sulbha R.Suple<sup>2</sup>

<sup>2</sup>Department of Mathematics,  
KarmavirDadasahebKannamwar  
College of Engineering,  
Nagpur 440009, India  
rajsulu4762@rediffmail.com

---

**Abstract:** *-In this paper,  $Z = (z - t)$  plane gravitational wave in higher dimensions is studied with the matter Wet Dark Energy in the framework of Rosen's Bimetric theory of gravitation [1] and observed that the Wet Dark Energy does not survive in this theory. Only a vacuum model can be constructed. Also some physical properties of the Models are discussed.*

**Keywords:** *Plane gravitational waves, Wet Dark Energy, Bimetric relativity.*

**AMS Code:** 83C05 (General Relativity)

---

### 1. INTRODUCTION

It is well known that the cosmological models based on General relativity contain an initial singular state (the big bang) from which the universe expands. The singular state can be avoided if the behavior of matter and radiation is described by the quantum theory. Unfortunately nobody has given a way to do this satisfactorily. A satisfactory physical theory should be free from singularities because the presence of singularity means a break-down of physical laws provided by the theory. Naturally, taking into consideration these singularities in general relativity one looks carefully at the foundation of general relativity and thinks whether modification can be made to improve it. With this Motivation, Rosen [1] proposed a bimetric theory of gravitation incorporating the covariance and equivalence principles. It is based on simple form of Lagrangian and has a simpler mathematical structure than that of the general theory of Relativity. In this theory at each point of space-time, there exist two metric tensors: a Riemannian metric tensor  $g_{ij}$  and the background flat space-time metric tensor  $\gamma_{ij}$ . The metric tensor  $g_{ij}$  determines the Riemannian geometry of the curved space time which plays the same role as given in the Einstein's general relativity and it interacts with matter. The background metric tensor  $\gamma_{ij}$  refers to the geometry of the empty (free from matter and radiation) universe and describes the inertial forces. This metric tensor  $\gamma_{ij}$  has no direct physical significance but appears in the field equations.

Therefore it interacts with  $g_{ij}$  but not directly with matter. One can regard  $\gamma_{ij}$  as giving the geometry that would exist if there were no matter. Moreover, the bimetric theory also satisfied the covariance and equivalence principles: the formation of general relativity. The theory agrees with the present observational facts pertaining to general relativity.

Thus at every point of space-time there are two line elements:

$$ds^2 = g_{ij} dx^i dx^j \quad (1.1)$$

And 
$$d\sigma^2 = \gamma_{ij} dx^i dx^j \quad (1.2)$$

Where  $ds$  is the interval between two neighbouring events as measured by means of a clock and measuring rod. The interval  $d\sigma$  is an abstract or geometrical quantity not directly measurable.

One can regard it as describing the geometry that would exist if no matter were present. Plane gravitational waves are usually discussed as a special case of the well-established plane fronted gravitational waves with parallel rays, the so called pp- waves. The method of specialization is quite technical, e.g. the curvature tensor must be complex recurrent with a recurrence vector which is collinear with a real null vector. H Takeno (1961) [3] propounded a rigorous discussion of plane gravitational waves, defined various terms by formulating a meaningful mathematical version and obtained numerous results.

A fairly general case of "plane" gravitational wave is represented by the metric

$$ds^2 = -A dx^2 - 2D dx dy - B dy^2 - dz^2 + dt^2 \quad (1.3)$$

both for weak field approximation and for exact solutions of Einstein field equations.

Reformulating Takeno's (1961) [3] definition of plane wave, we will use here,

$Z = z - t$  type plane gravitational waves by using the line elements,

$$ds^2 = -A dx^2 + dy^2 + du^2 - C dz^2 - dt^2 \quad (1.4)$$

Mohseni, Tucker and Wang [4] have studied the motion of spinning test particles in plane gravitational waves.

S Kessari, D Singh et al [5], analyzed the motion of electrically neutral massive spinning test particle in the plane gravitational and electromagnetic wave background. The theory of plane gravitational waves have been studied by many investigators Takeno [6]; Pandey [7]; Lal and Shafiullah [8]; Lu Huiqing [9]; Bondi, H.et.al.[10],Torre,C.G.[11]; Hogan, P.A.[13]; and they obtained the solutions .

Further, Sahoo P.K. [14] have studied Inhomogeneous plane symmetric string cosmological models in Bimetric theory, Bhoyar and Deshmukh [15] have studied  $Z = (z-t)$  type plane fronted waves and electromagnetic waves with massless scalar plane wave and massive scalar plane waves in Peres space-time. Also Sahoo, Behera, Tripathy et al [16] have investigated inhomogeneous cosmological models in Bimetric theory of Gravitation.

Also Deo and Ronghe [17],[18] ; Deo and Suple[20] ,[21]have studied plane gravitational waves and obtained the solutions.

In this paper, we will study  $Z = z - t$  type plane gravitational wave with wet dark energy and will observe the result in the context of Bimetric theory of relativity.

## 2. FIELD EQUATIONS IN BIMETRIC RELATIVITY

Rosen N. has proposed the field equations of Bimetric Relativity from variation principle as

$$K_i^j = N_i^j - \frac{1}{2} N g_i^j = -8\pi\kappa T_i^j \quad (2.1)$$

$$\text{Where } N_i^j = \frac{1}{2} \gamma^{\alpha\beta} [g^{hj} g_{hi} |_{\alpha}] |_{\beta} \quad (2.2)$$

$$N = N_{\alpha}^{\alpha} \quad \kappa = \sqrt{\frac{g}{\gamma}} \quad (2.3)$$

$$\text{and } g = |g_{ij}| \quad , \gamma = |\gamma_{ij}| \quad (2.4)$$

Where a vertical bar (|) denotes a covariant differentiation with respect to  $\gamma_{ij}$

And,  $T_i^j$  the energy momentum tensor for wet dark energy is given by

$$T_i^j = T_{i\ wdf}^j = \rho_{wdf} + p_{wdf} u_i u^j - p_{wdf} g_i^j \quad (2.5)$$

together with  $g_i^j u_i u^j = 1, u_5 u^5 = 1$  where  $u_i$  is the five-velocity vector of the fluid having  $p$  and  $\rho$  as proper pressure and energy density respectively.

In co-moving coordinate system we have

$$T_1^1 = T_2^2 = T_3^3 = T_4^4 = -p_{wdf}, T_5^5 = \rho_{wdf} \text{ and } T_i^j = 0 \text{ for } i \neq j$$

### 3. Z= z-t TYPE HIGHER FIVE DIMENSIONAL PLANE GRAVITATIONAL WAVE WITH WET DARK ENERGY

For plane gravitational wave  $Z = z - t$ , we have the line element as,

$$ds^2 = -A(dx^2 + dy^2 + du^2) - C(dz^2 - dt^2) \quad (3.1)$$

Where  $A = A(Z), C = C(Z)$  and  $Z = z - t$

Corresponding to equation (3.1), we consider the line element for background metric  $\gamma_{ij}$  as

$$d\sigma^2 = -(dx^2 + dy^2 + du^2 + dz^2) + dt^2 \quad (3.2)$$

Since  $\gamma_{ij}$  is the Lorentz metric i.e. (-1,-1,-1,-1, 1), therefore  $\gamma$ -covariant derivative becomes the ordinary partial derivative.

Using equations (2.1) to (2.5) with (3.1) and (3.2), we get,

$$\frac{1}{4} \left[ \left( \frac{A'^2}{A^2} - \frac{A''}{A} \right) - \left( \frac{\dot{A}^2}{A^2} - \frac{\ddot{A}}{A} \right) \right] + \frac{1}{2} \left[ \left( \frac{C'^2}{C^2} - \frac{C''}{C} \right) - \left( \frac{\dot{C}^2}{C^2} - \frac{\ddot{C}}{C} \right) \right] = -16\pi\kappa p_{wdf} \quad (3.3)$$

$$\frac{3}{4} \left[ \left( \frac{A'^2}{A^2} - \frac{A''}{A} \right) - \left( \frac{\dot{A}^2}{A^2} - \frac{\ddot{A}}{A} \right) \right] = -16\pi\kappa p_{wdf} \quad (3.4)$$

$$\frac{3}{4} \left[ \left( \frac{A'^2}{A^2} - \frac{A''}{A} \right) - \left( \frac{\dot{A}^2}{A^2} - \frac{\ddot{A}}{A} \right) \right] = 16\pi\kappa\rho_{wdf} \quad (3.5)$$

where the overhead primes and dots denotes differentiation of the metric potentials with respect to  $z$  and  $t$  resp.

To solve the field equations (3.3) – (3.5), we note that there are three equations connecting four unknowns,  $A, C, p_{wdf}$  and  $\rho_{wdf}$ . So one relation connecting these variables is needed.

Here assume the relation between the metric potentials such as Bhattacharya and Karade [12]

$$C = \alpha A, \quad \alpha \neq 0 \text{ is a constant.} \quad (3.6)$$

on solving (3.3) to (3.5) we obtain,

$$p_{wdf} + \rho_{wdf} = 0 \quad (3.7)$$

In view of the reality conditions i.e.  $p > 0, \rho > 0$  must hold.

The above conditions (3.6) is satisfied only when

$$p = 0, \rho = 0. \quad (3.8)$$

This means that the physical parameters, viz proper pressure ( $p$ ), energy density ( $\rho$ ) are identically zero.

Thus higher five dimensional plane gravitational wave with wet dark energy in biometric relativity does not survive and hence only vacuum model is obtained.

Using (3.7), the vacuum field equations are

$$\frac{1}{4} \left[ \left( \frac{A'^2}{A^2} - \frac{A''}{A} \right) - \left( \frac{\dot{A}^2}{A^2} - \frac{\ddot{A}}{A} \right) \right] + \frac{1}{2} \left[ \left( \frac{C'^2}{C^2} - \frac{C''}{C} \right) - \left( \frac{\dot{C}^2}{C^2} - \frac{\ddot{C}}{C} \right) \right] = 0 \quad (3.9)$$

$$\frac{3}{4} \left[ \left( \frac{A'^2}{A^2} - \frac{A''}{A} \right) - \left( \frac{\dot{A}^2}{A^2} - \frac{\ddot{A}}{A} \right) \right] = 0 \quad (3.10)$$

Thus field equations yields

$$\left[ \left( \frac{A'^2}{A^2} - \frac{A''}{A} \right) - \left( \frac{\dot{A}^2}{A^2} - \frac{\ddot{A}}{A} \right) \right] = 0 \quad (3.11)$$

By using Method of separation of Variables, (3.11) gives us the solution

$$A Z = e^{\left[ \frac{l}{2} z^2 + t^2 + l_1 z + l_2 t \right]} \quad (3.12)$$

With the help of (3.6) we get

$$C Z = \alpha e^{\left[ \frac{l}{2} z^2 + t^2 + l_1 z + l_2 t \right]} \quad (3.13)$$

where,  $l, l_1$  and  $l_2$  are the constants of integration.

Thus substituting the value of (3.12) and (3.13) in (3.1), we get the vacuum line element as

$$ds^2 = - \mathbf{exp} \left[ \frac{l}{2} z^2 + t^2 + l_1 z + l_2 t \right] dx^2 + dy^2 + du^2 - \alpha \mathbf{exp} \left[ \frac{l}{2} z^2 + t^2 + l_1 z + l_2 t \right] dz^2 - dt^2 \quad (3.14)$$

The model (3.14) has no singularities either at  $z = 0$  or at  $t = 0$  and it does have singularities as  $z \rightarrow \infty$  and  $t \rightarrow \infty$

#### 4. CONCLUSION

In the study of  $Z = z - t$  type plane gravitational waves, there is nil contribution of wet dark energy in Bimetric theory of relativity respectively. It is observed that the matter fieldlike wetdark energy cannot be a source of gravitational field in the Rosen's bimetric theory but only vacuum model exists.

#### ACKNOWLEDGEMENT

The authors are thankful to Dr. R. D. Giri, Prof. Emeritus, P.G.T.D. (Mathematics), R. T. M. N. U., Nagpur, India for his constant inspiration.

**REFERENCES**

- [1]. Rosen N. (1940) General Relativity and Flat space I. Phys. Rev.57, 147.
- [2]. Rosen N. (1973) A bimetric theory of gravitation I Gen. Relat.Grav.04, 435-47.
- [3]. Takeno H. (1961) The mathematical theory of plane gravitational waves in General Relativity. Scientific report of Research Institute for theoretical Physics, Hiroshima University, Hiroshima, Ken, Japan (1961).
- [4]. Mohseni, M.; Tucker, R. W.; Wang, C. (2001). On the motion of spinning test particles in Plane gravitational waves. Class. Quantum Grav. 18 3007-3017
- [5]. Kessari, S ; Singh, D. et al, (2002). Scattering of spinning test particles by plane gravitational and Electromagnetic waves.gr-qc/0203038,Class.Quant.Grav. 194943-4952
- [6]. Takeno, H. (1958). A comparison of plane wave solutions in general relativity with those in non-symmetric theory. Prog.Theor.Phys.20, 267-276
- [7]. Pandey, S. N. (1979). Plane wave solutions in Finzi's non-symmetric unified field theory. Theo. Math. Phys. 39, 371-375
- [8]. Lal, K. B. ; Shafiullah, (1980). On plane wave solutions of non – symmetric field equations of unified theories of Einstein Bonner and Schrödinger. Annali de Mathematica ed pura Applicata.,126, 285-298.
- [9]. Lu Huiqing (1988). Plane gravitational waves under a non - zero cosmological constant. hin. Astronomy Astrophys. 12, 186-190.
- [10]. Bondi, H. ; Pirani, F.A.E. and Robinson, I.(1959). Gravitational waves in general relativity III. Exact plane waves. Proc. Roy.Soc.Lond.A23, 25, 519-533
- [11]. Torre, C.G. (2006) Gravitational waves – Just plane symmetry. Gen. Relat. Grav. 38, 653-662
- [12]. Bhattacharya and Karade T.M.(1993) Astrophys.SpaceSci.202,69
- [13]. Hogan, P.A.(1999). Gravitational waves and Bertotti-Robinson space- time. Math. Proc. Roy. Irish Acad. 99A, 51-55.
- [14]. Sahoo P.K. (2009) Int J TheorPhyd 48; 2022-2026 DOI10.1007/s10773-009-9977-1
- [15]. Bhoyar S. R, Deshmukh A.G Romanian Reports in Physics, Vol.63, No.1, P.25-34, 2011
- [16]. Sahoo, Behera, Tripathy et al Communication in Physics, Vol.20, No.2, pp.121-127
- [17]. Ronghe A.K. and Deo S.D. - Plane Gravitational Waves In Bimetric Relativity. JVR (2011) 6, 1-11
- [18]. Ronghe A.K. and Deo S.D. Plane gravitational waves with wet dark energy. International Journal of Mathematical Archive- 2(3) Mar.-391-392
- [19]. Deo S.D. and Suple S.R. Plane gravitational waves with cosmic strings in Bimetric Relativity. Mathematica Aeterna, Vol. 3, 2013, no. 6, 489 - 496
- [20]. Deo S.D. and Suple S.R. Plane gravitational waves with Macro and Micro Matter fields in Bimetric Relativity. International Journal of Mathematics Trends and Technology – Volume 8 Number 1 – April 2014.