

Another Look at Planetary Motion

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Abstract: *In this brief note, the classical and relativistic equations of planetary motion are studied using polar coordinates in a different way than the standard approach. Conditions are then given when the relativistic equations yield bounded solutions.*

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1. INTRODUCTION

In this note, a straightforward account will be given of the well-known differential equations of planetary motion. However, we will discuss them using the (r, θ) coordinates in a slightly different way than the usual case (see [1, pp. 231-235] for the classical approach and [2, pp. 471-496] for a thorough presentation of the behavior of a body in motion subject to a central force). This discussion simplifies the analysis of these fundamental equations of celestial mechanics (see [3] for an excellent introduction to this field).

2. METHODS

The first equation is the classical Newtonian model while the second one is derived from general relativity (see [4, pp. 270-276] for further details). The two equations are:

(1) $u''(\theta) + u(\theta) = c$ (where $u = 1/r$, r being the radius from the given object to one of the foci which is the sun for the purpose of this discussion and c is a certain positive constant) and

(2) $u''(\theta) + u(\theta) - c_1 u(\theta)^2 = c_2$ (where c_1 and c_2 are other positive constants).

In our discussion, we shall solve the first equation exactly and the second one implicitly by treating it as a two dimensional system. This will give us some clarity into the behavior of their solutions.

3. MAIN RESULTS AND DISCUSSION

We first begin our analysis by solving equation (1) which has the general solution $u = c + a \sin(\theta) + b \cos(\theta)$ where a and b are constants. Noting that $u = 1/r$ we may therefore conclude that $r = (c + a \sin(\theta) + b \cos(\theta))^{-1}$ which is the general form of an ellipse, the planetary

orbit. This general solution covers all cases including the case when the major and minor axes of the ellipse are not parallel or perpendicular to the XY axes. When setting $a = 0$, we have the equation of the ellipse in standard form $r = 1/(c + b \cos(\theta))$. The only constraint is $c + b \cos(\theta) \neq 0$.

As far as equation (2) is concerned we can show directly that all solutions are bounded when given certain initial conditions. First, multiply (2) by $2u'$ and integrate for 0 to θ obtaining

$$(3) \quad u'(\theta)^2 + u(\theta)^2 - 2c_1 u(\theta)^3 / 3 = 2c_2 u(\theta) - 2c_2 u(0) + u'(0)^2 + u(0)^2 - 2c_1 u(0)^3 / 3.$$

Next, using the fact that $u = 1/r$ and $u' = -r'/r^2$ and then multiplying equation (3) by r^4 transforms equation (3) into

$$(4) \quad r'(\theta)^2 + r(\theta)^2 - 2c_1 r(\theta) / 3 = 2c_2 r(\theta)^3 - 2c_2 r(\theta)^4 u(0) + k r(\theta)^4$$

where $k = u'(0)^2 + u(0)^2 - 2c_1 u(0)^3 / 3$. If $k - 2c_2 u(0) < 0$, then should $r \rightarrow \infty$ the LHS of (4) approaches ∞ while the RHS approaches $-\infty$ which is impossible. In other words, the solutions

must remain bounded as $t \rightarrow \infty$ given these conditions. Should $k - 2c_2u(0) \geq 0$, then the solutions may be unbounded (note: the equilibrium points of (2) in the (u, θ) plane correspond to regions of unboundedness in the (r, θ) plane).

We could also look at equation (2) another way by transforming it into the following dynamical system

$$(5) \quad \begin{aligned} u' &= v \\ v' &= -u + c_1u^2 + c_2. \end{aligned}$$

Then, we convert (5) into the first order differential equation,

$$(6) \quad dv/du = (-u + c_1u^2 + c_2)/v. \text{ This becomes}$$

$$(7) \quad v \, dv = (-u + c_1u^2 + c_2) \, du.$$

Finally, integrating equation (7) from 0 to θ yields the following result

$$(8) \quad \frac{1}{2}(v(\theta)^2 - v(0)^2) = -\frac{u(\theta)^2}{2} + \frac{c_1u(\theta)^3}{3} + c_2u(\theta) - \frac{u(0)^2}{2} + \frac{c_1u(0)^3}{3} - c_2u(0).$$

Letting $u = 1/r$ and $u' = v = -r'/r^2$ again yield equation (4) after some algebraic manipulation.

4. CONCLUSION

Using standard methods from differential equations, the above analysis clearly gives a straightforward and precise analysis of planetary motion both classical and relativistic which plays an important role in celestial mechanics.

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