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Gauss Symbol and Unification of Gravity and Electromagnetic

Field in Reissner-Nordstrom and Kerr-Newman Solution

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Abstract: Solutions of unified theory equations of gravity and electromagnetism satisfy Einstein-Maxwell equation. Hence, solutions of the unified theory is Reissner-Nordstrom solution, Kerr-Newman solution in vacuum. In this careful point, Einstein gravity field equation's energy-momentum tensor has two sign. Therefore, according to the solution, the sign is treated. We found in revised Einstein gravity tensor equation, the condition is satisfied by 2-order contravariant metric tensor two times product and Gauss symbol.

Keywords: General relativity theory, Unified Theory; 2-order contravariant metric tensor two times product; Gauss Symbol Energy-momentum's sign

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1. Introduction

This theory's aim is that we discover the revised Einstein gravity equation had Reissner-Nordstrom solution and Kerr-Newman solution in vacuum. In this careful point, Einstein gravity field equation's energy-momentum tensor has two sign. Therefore, according to the solution, the sign is treated.

First, new symmetric revised Einstein gravity field Equation is

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R + \Lambda g_{\mu\nu} (-1)^{\left[\frac{\mu+\nu}{4}\right]} (g^{\theta\theta})^2 = \pm \frac{8\pi G}{C^4} T_{\mu\nu}$$
 (1)

Gauss symbol [x] is the greatest integer that is less than or equal to x. For example,

$$[3.141592] = 3, [2.71828] = 2, [-0.802] = -1,...$$
 (2)

Hence, new term is in Kerr-Newman solution,

$$(\mu, \nu) = (0, 0) \rightarrow (-1)^{[\frac{\mu+\nu}{4}]} = 1 , (\mu, \nu) = (1, 1) \rightarrow (-1)^{[\frac{\mu+\nu}{4}]} = 1$$

$$(\mu, \nu) = (0, 3) \to (-1)^{[\frac{\mu + \nu}{4}]} = 1 \ (\mu, \nu) = (3, 0) \to (-1)^{[\frac{\mu + \nu}{4}]} = 1$$

$$(\mu, \nu) = (2, 2) \rightarrow (-1)^{\left[\frac{\mu+\nu}{4}\right]} = -1 , (\mu, \nu) = (3, 3) \rightarrow (-1)^{\left[\frac{\mu+\nu}{4}\right]} = -1$$
 (3)

We take covariant differential operator,

$$(R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R)_{;\mu} + \Lambda g_{\mu\nu} (-1)^{[\frac{\mu+\nu}{4}]} 2g^{\theta\theta}g^{\theta\theta}_{;\mu} = \pm \frac{8\pi G}{C^4}T_{\mu\nu;\mu} = 0$$
(4)

$$(R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R)_{\nu} + \Lambda g_{\mu\nu} (-1)^{[\frac{\mu+\nu}{4}]} 2g^{\theta\theta}g^{\theta\theta}_{\nu} = \pm \frac{8\pi G}{G^4}T_{\mu\nu\nu} = 0$$
 (5)

In this time, in Kerr-Newman solution [4]

$$g_{\theta\theta} = 1/g^{\theta\theta} = \rho^2 = r^2 + a^2 \cos^2 \theta \tag{6}$$

$$g^{\theta\theta}_{,\rho} = \frac{\partial g^{\theta\theta}}{\partial x^{\rho}} + 2\Gamma^{\theta}_{\sigma\rho}g^{\sigma\theta} = \frac{\partial g^{\theta\theta}}{\partial r} + 2\Gamma^{\theta}_{\theta r}g^{\theta\theta}$$

$$= \frac{\partial}{\partial r} \left(\frac{1}{\rho^{2}}\right) + 2 \cdot \frac{r}{\rho^{2}} \cdot \frac{1}{\rho^{2}} = -2\frac{1}{\rho^{3}} \cdot \frac{r}{\rho} + \frac{2r}{\rho^{4}} = 0$$

$$g^{\theta\theta}_{,\rho} = \frac{\partial g^{\theta\theta}}{\partial x^{\rho}} + 2\Gamma^{\theta}_{\sigma\rho}g^{\sigma\theta} = \frac{\partial g^{\theta\theta}}{\partial \theta} + 2\Gamma^{\theta}_{\theta\theta}g^{\theta\theta}$$

$$= \frac{\partial}{\partial \theta} \left(\frac{1}{\rho^{2}}\right) - 2 \cdot \frac{1}{\rho^{2}} \cdot \frac{1}{2}\rho \frac{2a^{2}}{\rho} \cos\theta \sin\theta \cdot \frac{1}{\rho^{2}}$$

$$= -2 \cdot \frac{1}{\rho^{3}} \cdot -\frac{2a^{2}\cos\theta \sin\theta}{\rho} - \frac{4a^{2}}{\rho^{4}} \cos\theta \sin\theta = 0$$
(8)

$$0 = V_{\rho} = \frac{\partial_{X} {}^{\rho}}{\partial_{X} {}^{\rho}} V_{\alpha}^{\dagger}, \quad V_{\alpha}^{\dagger} = 0$$

$$(9)$$

Therefore, if the coordinate is not the spherical coordinate, the covariant differential of $g^{\theta\theta} = \frac{1}{\rho^2}$ is

still zero in the changed coordinate

2. THE NEW REVISED EINSTEIN GRAVITY EQUATION AND REISSNER-NORDSTROM SOLUTION

In this theory, Eq(1) can change the following equation. This time, energy-momentum tensor's sign is minus in Einstein equation in Reissner-Nordstrom space-time,

$$R_{\mu\nu} = -\frac{8\pi G}{G^4} (\Gamma_{\mu\nu} - \frac{1}{2} g_{\mu\nu} T^{\lambda}_{\lambda}) + \Lambda g_{\mu\nu} (-1)^{[\frac{\mu+\nu}{4}]} (g^{\theta\theta})^2$$
 (10)

In this time, in vacuum, Eq(10) is in Reissner-Nordstrom solution.

$$R_{\mu\nu} = \Lambda g_{\mu\nu} \left(-1\right)^{\left[\frac{\mu+\nu}{4}\right]} \left(g^{\theta\theta}\right)^{2} = \Lambda g_{\mu\nu} \left(-1\right)^{\left[\frac{\mu+\nu}{4}\right]} \frac{1}{r^{4}}$$
(11)

Reissner-Nordstrom solution of Einstein-Maxwell equation is

$$g_{00} = -1 + \frac{2GM}{m^2} - \frac{kGQ^2}{r^2C^4}, g_{11} = 1/\left(1 - \frac{2GM}{m^2} + \frac{kGQ^2}{r^2C^4}\right)$$

$$g_{22} = r^2, g_{33} = r^2 \sin^2 \theta$$
(12)

The proper time of spherical coordinates is

$$d\tau^{2} = A(t,r)dt^{2} - \frac{1}{c^{2}}[B(t,r)dr^{2} + r^{2}d\theta^{2} + r^{2}\sin^{2}\theta d\varphi^{2}]$$

$$(\mu, \nu) = (0, 0) \rightarrow (-1)^{\left[\frac{\mu+\nu}{4}\right]} = 1$$
, $(\mu, \nu) = (1, 1) \rightarrow (-1)^{\left[\frac{\mu+\nu}{4}\right]} = 1$

$$(\mu, \nu) = (2, 2) \rightarrow (-1)^{\left[\frac{\mu+\nu}{4}\right]} = -1, (\mu, \nu) = (3, 3) \rightarrow (-1)^{\left[\frac{\mu+\nu}{4}\right]} = -1$$
 (13)

If we use Eq(11), Eq(13), we obtain the Ricci-tensor equations. [3]

$$R_{tt} = -\frac{A''}{2B} + \frac{A'B'}{4B^2} - \frac{A'}{Br} + \frac{A^{\ell}}{4AB} + \frac{\ddot{B}}{2B} - \frac{\dot{B}^2}{4B^2} - \frac{\dot{A}\dot{B}}{4AB} = -\Lambda A \frac{1}{r^4}$$
 (14)

$$R_{x} = \frac{A''}{2A} - \frac{A^{2}}{4A^{2}} - \frac{A'B'}{4AB} - \frac{B'}{Bx} - \frac{B'}{2A} + \frac{\dot{A}\dot{B}}{4A^{2}} + \frac{\dot{B}^{2}}{4AB} = \Lambda B \frac{1}{x^{4}}$$
(15)

$$R_{\theta\theta} = -1 + \frac{1}{B} - \frac{rB'}{2B^2} + \frac{rA'}{2AB} = -\Lambda r^2 \frac{1}{r^4} = -\Lambda \frac{1}{r^2}$$
 (16)

$$R_{\theta\theta} = R_{\theta\theta} \sin^2 \theta \tag{17}$$

$$R_{tr} = -\frac{\dot{B}}{Br} = 0 \tag{18}$$

$$R_{t\theta} = R_{t\theta} = R_{x\theta} = R_{x\theta} = R_{\theta\theta} = 0 \tag{19}$$

In this time, $'=\frac{\partial}{\partial r}$, $\cdot=\frac{1}{c}\frac{\partial}{\partial t}$

If we calculate,

$$\frac{R_{\pm}}{A} + \frac{R_{\pm}}{B} = -\frac{1}{Br} \left(\frac{A'}{A} + \frac{B'}{B} \right) = -\frac{(AB)'}{rAB^2} = 0$$
 (20)

Hence, we obtain this result.

$$A = \frac{1}{B} \tag{21}$$

If Eq(21) inserts Eq(16),

$$R_{\theta\theta} = -1 + \frac{1}{B} - \frac{rB'}{2B^2} + \frac{rA'}{2AB} = -1 + (\frac{r}{B})' = -\Lambda \frac{1}{r^2}$$
 (22)

If we solve Eq(22),

$$\frac{r}{B} = r + C + \frac{\Lambda}{r}$$

$$\rightarrow A = \frac{1}{B} = 1 - \frac{2GM}{rc^2} + \frac{kGQ^2}{r^2c^4}, \quad \Lambda = k\frac{GQ^2}{c^4}, C = -\frac{2GM}{c^2}$$
 (23)

Therefore, in vacuum, the spherical solution of the revised Einstein gravity equation is Reissner-Nordstrom solution.

$$d\tau^{2} = \left(1 - \frac{2GM}{rc^{2}} + \frac{kGQ^{2}}{r^{2}c^{4}}\right)dt^{2} - \frac{1}{c^{2}}\left[\frac{dr^{2}}{\left(1 - \frac{2GM}{rc^{2}} + \frac{kGQ^{2}}{r^{2}c^{4}}\right)} + r^{2}d\theta^{2} + r^{2}\sin^{2}\theta d\phi^{2}\right]$$
(24)

3. KERR-NEWMAN SOLUTION AND NEW REVISED GRAVITY FIELD EQUATION

In this theory, Eq(1) can change the following equation. In this time, energy-momentum tensor's sign is plus in Einstein equation in Kerr-Newman space-time,

In this time, in vacuum, Eq(25) is in Kerr-Newman solution,

$$R_{\mu\nu} = \Lambda g_{\mu\nu} (-1)^{\left[\frac{\mu+\nu}{4}\right]} (g^{\theta\theta})^2 = \Lambda g_{\mu\nu} (-1)^{\left[\frac{\mu+\nu}{4}\right]} \frac{1}{\rho^4}$$
 (26)

We think Kerr-Newman equation and solution.

Kerr-Newman solution is

$$ds^{2} = g_{00}c^{2}dt^{2} + 2g_{03}cdtd\phi + g_{11}dr^{2} + g_{22}d\theta^{2} + g_{33}d\phi^{2}$$

$$(\mu, \nu) = (0, 0) \rightarrow (-1)^{\left[\frac{\mu+\nu}{4}\right]} = 1$$
, $(\mu, \nu) = (1, 1) \rightarrow (-1)^{\left[\frac{\mu+\nu}{4}\right]} = 1$

$$(\mu, \nu) = (0, 3) \rightarrow (-1)^{\left[\frac{\mu+\nu}{4}\right]} = 1 \ (\mu, \nu) = (3, 0) \rightarrow (-1)^{\left[\frac{\mu+\nu}{4}\right]} = 1$$

$$(\mu, \nu) = (2, 2) \to (-1)^{[\frac{\mu + \nu}{4}]} = -1 \ , (\mu, \nu) = (3, 3) \to (-1)^{[\frac{\mu + \nu}{4}]} = -1$$

$$\rho^2 = r^2 + a^2 \cos^2 \theta \tag{27}$$

Hence, Kerr-Newman equation is [5]

$$R_{tt} = \Lambda g_{00} \frac{1}{\rho^4} \tag{28}$$

$$R_{x} = \Lambda g_{11} \frac{1}{\rho^4} \tag{29}$$

$$R_{\theta\theta} = -\Lambda g_{22} \frac{1}{\rho^4} = -\Lambda \frac{1}{\rho^2} \tag{30}$$

$$R_{\phi\phi} = -\Lambda g_{33} \frac{1}{\rho^4} \tag{31}$$

$$R_{t\phi} = R_{\phi t} = \Lambda g_{03} \frac{1}{\rho^4} = \Lambda g_{30} \frac{1}{\rho^4}$$
 (32)

$$R_{\phi t} = R_{t\phi} = \Lambda g_{03} \frac{1}{\rho^4} = \Lambda g_{30} \frac{1}{\rho^4}$$
 (33)

Otherwise $R_{\mu\nu} = 0$

In this time, for proving Eq(28)-Eq(33), Maxwell tensor equations of Einstein-Maxwell equation are

$$-\Lambda g_{\mu\nu} (-1)^{\left[\frac{\mu+\nu}{4}\right]} (g^{\theta\theta})^{2} = R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = \frac{8\pi G}{c^{4}} T_{\mu\nu}$$

$$= \frac{2G}{c^{5}} (F_{\mu\rho} F_{\nu}^{\ \rho} - \frac{1}{4} g_{\mu\nu} F_{\rho\sigma} F^{\rho\sigma})$$

$$= \frac{2G}{c^{5}} (g_{\mu\nu} F_{\nu\rho} F^{\nu\rho} - \frac{1}{4} g_{\mu\nu} F_{\rho\sigma} F^{\rho\sigma})$$
(34)

In Kerr-Newman coordinates, if we calculate Eq(34) of g_{03} , g_{30} by g_{30} [4,7]

$$\frac{8\pi G}{c^4} T_{\mu\nu} = \frac{2G}{c^5} \left(g_{\mu\nu} F_{\nu\rho} F^{\nu\rho} - \frac{1}{4} g_{\mu\nu} F_{\rho\sigma} F^{\rho\sigma} \right)
= \frac{8\pi G}{c^4} T_{30} = \frac{8\pi G}{c^4} T_{03} = \frac{2G}{c^5} \left(g_{30} F_{01} F^{01} - \frac{1}{4} g_{30} Q B^2 - 2E^2 \right) \right)
= \frac{2G}{c^5} \left(-g_{30} E^2 - \frac{1}{2} g_{30} B^2 - E^2 \right) = -\frac{G}{c^5} B^2 + E^2 g_{30} = -k \frac{G}{c^4} \frac{Q^2}{\rho^4} g_{30} \right)
= -\Lambda g_{30} \left(-1 \right)^{\frac{(0+3)}{4}} \frac{1}{\rho^4} = -\Lambda g_{30} \frac{1}{\rho^4} = -\Lambda g_{03} \frac{1}{\rho^4} = -g_{03} \left(-1 \right)^{\frac{(0+3)}{4}} \frac{1}{\rho^4}$$

$$E = F_{01} = -F_{10} \frac{Q(r^2 - d \cos \theta)}{\rho^4}, \quad B = F_{23} = -F_{32} = \frac{2Qar \cos \theta}{\rho^4}
B^2 + E^2 = \frac{Q^2}{\rho^4} \tag{35}$$

In this careful point, if we calculate Eq(34) by g_{03} , the calculation is wrong.

In Kerr-Newman coordinates, if we calculate Eq(34) of g_{00} [4,7]

$$\frac{2G}{c^{5}} \left(g_{\mu\nu} F_{\nu\rho} F^{\nu\rho} - \frac{1}{4} g_{\mu\nu} F_{\rho\sigma} F^{\rho\sigma} \right)
= \frac{2G}{c^{5}} \left(g_{00} F_{01} F^{01} - \frac{1}{4} g_{00} Q B^{2} - 2 E^{2} \right) \right)
= \frac{2G}{c^{5}} \left(-g_{00} E^{2} - \frac{1}{2} g_{00} \mathcal{B}^{2} - E^{2} \right) \right) = -\frac{G}{c^{5}} \mathcal{B}^{2} + E^{2} \right) g_{00} = -k \frac{G}{c^{4}} \frac{Q^{2}}{\rho^{4}} g_{00}
= -\Lambda g_{00} \left(-1 \right)^{\frac{O}{4}} \frac{1}{\rho^{4}} = -\Lambda g_{00} \frac{1}{\rho^{4}}
E = F_{01} = -F_{10} \stackrel{Q}{=} \frac{(r^{2} - d \cos \theta)}{\rho^{4}}, \quad B = F_{23} = -F_{32} = \frac{2Qar \cos \theta}{\rho^{4}}
B^{2} + E^{2} = \frac{Q^{2}}{\rho^{4}}$$
(36)

In Kerr-Newman coordinates, if we calculate Eq(34) of $g_{11}[4,7]$

$$\begin{split} &\frac{2G}{c^{5}} \left(g_{\mu\nu} F_{\nu\rho} F^{\nu\rho} - \frac{1}{4} g_{\mu\nu} F_{\rho\sigma} F^{\rho\sigma} \right) \\ &= \frac{2G}{c^{5}} \left(g_{11} F_{10} F^{10} - \frac{1}{4} g_{11} \left(2B^{2} - 2E^{2} \right) \right) \\ &= \frac{2G}{c^{5}} \left(-g_{11} E^{2} - \frac{1}{2} g_{11} \left(B^{2} - E^{2} \right) \right) = -\frac{G}{c^{5}} \left(B^{2} + E^{2} \right) g_{11} = -k \frac{G}{c^{4}} \frac{Q^{2}}{\rho^{4}} g_{11} \\ &= -\Lambda g_{11} \left(-1 \right)^{\left[\frac{1+1}{4} \right]} \frac{1}{\rho^{4}} = -\Lambda g_{11} \frac{1}{\rho^{4}} \end{split}$$

$$E = F_{01} = -F_{10} = \frac{Q(r^2 - a^2 \cos^2 \theta)}{\rho^4}, \quad B = F_{23} = -F_{32} = \frac{2Qar \cos \theta}{\rho^4}$$

$$B^2 + E^2 = \frac{Q^2}{\rho^4}$$
(37)

In Kerr-Newman coordinates, if we calculate Eq(34) of g_{22} [4,7]

$$\frac{2G}{c^{5}} \langle g_{\mu} F \rangle_{\rho}^{\nu} \rho - \frac{1}{4} g \rangle_{\mu} F_{\rho}^{\rho} \rho^{\rho} \rho^{\rho}$$

In Kerr-Newman coordinates, if we calculate Eq(34) of g_{33} [4,7]

$$\frac{2G}{c^{5}} (g_{\mu} F F_{\rho}^{\nu \rho} - \frac{1}{4} g F_{\mu \nu} F_{\rho \sigma}^{\rho})^{5}$$

$$= \frac{2G}{c^{5}} (g_{33} F_{32} F^{32} - \frac{1}{4} g_{33} (2B^{2} - 2E^{2}))$$

$$= \frac{2G}{c^{5}} (g_{33} B^{2} - \frac{1}{2} g_{33} (B^{2} - E^{2})) = \frac{G}{c^{5}} (B^{2} + E^{2}) g_{33} = k \frac{G}{c^{4}} \frac{Q^{2}}{\rho^{4}} g_{33}$$

$$= -\Lambda g_{33} (-1)^{\left[\frac{3+3}{4}\right]} \frac{1}{\rho^{4}} = \Lambda g_{33} \frac{1}{\rho^{4}}$$

$$E = F_{01} = -F_{10} = \frac{Q(r^{2} - d^{2} \cos \theta)}{\rho^{4}}, \quad B = F_{23} = -F_{32} = \frac{2Qar \cos \theta}{\rho^{4}}$$

$$B^{2} + E^{2} = \frac{Q^{2}}{\rho^{4}}$$
(39)

According to Eq(35-39), in Kerr-Newman solution, the electromagnetic energy-momentum tensor's formula is equal with the new revised Einstein equation's added term of our unified theory. Hence, we know if we solve the revised Einstein's equation, we can obtain Kerr-Newman solution.

4. CONCLUSION

We found the revised Einstein equation of unified theory (the gravity and electromagnetic field). This theory's strong point is 4-dimensional theory. This theory is different from 5-dimensional Kaluza-Klein theory. But as the method of describing universe, Einstein normal gravity equation is equal with

the revised Einstein equation of the unified theory because the electric charge has to be zero. .

REFERENCES

[1]S.Yi, "Spherical Solution of Classical Quantum Gravity", International Journal of Advanced Research in Physical Science, 6,8,2019, pp3-6

[2]S.Yi, "Unified Theory of Gravity and Electromagnetic Field", International Journal of Advanced Research in Physical Science, 6,10,2019,pp10-12

[3]S.Yi, "Unified Theory of Gravity and Electromagnetic Field in Reissner-Nodstrom, Kerr-Newman Solution", International Journal of Advanced Research in Physical Science, 6,11,2019, pp1-3

[4]S.Yi, "Unification for Gravity and Electromagnetic Field in Kerr-Newman Solution", International Journal of Advanced Research in Physical Science, 6,12,2019, pp11-13

[5]S.Yi, "Unified Theory of Gravity and Electromagnetic Field with Problem in Reissner-Nodstrom and Kerr-Newman Solution", International Journal of Advanced Research in Physical Science, 7,1,2020, pp12-15

[6]S.Yi, "Gauss Symbol and Unified Theory of Gravity and electromagnetic Field", International Journal of Advanced Research in Physical Science, 7, 7,2020, pp8-13,

[7]K.Rosquist, "Interacting Kerr-Newman Fields", International Center for Relativistic Astrophysics, MG12, Paris (2009)

[8]S. Weinberg, Gravitation and Cosmology(John wiley & Sons, Inc, 1972)

[9]P.Bergman, Introduction to the Theory of Relativity(Dover Pub. Co., Inc., New York, 1976), Chapter V

[10]C.Misner, K,Thorne and J. Wheeler, Gravitation(W.H.Freedman & Co.,1973)

[11]S.Hawking and G. Ellis, The Large Scale Structure of Space-Time(Cam-bridge University Press, 1973)

[12]R.Adler,M.Bazin and M.Schiffer,Introduction to General Relativity(McGraw-Hill,Inc.,1965)

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